

1) Electrons in the CB of a nondegenerate semiconductor

a) Let's consider the energy distribution of electrons $n_E(E)$ in the conduction band (CB). Assuming that the density of state

$$g_{CB}(E) \propto (E - E_c)^{1/2}$$

and using Boltzmann statistics

$$f(E) \approx e^{-\frac{E - E_F}{k_B T}}$$

we show here that the energy distribution of the electrons in the CB can be written as

$$n_x(x) = C \sqrt{x} e^{-x}$$

where

$$x = \frac{E - E_c}{k_B T}$$

is the electron energy in terms of $k_B T$ measured from E_c and C is a constant at a given temperature (independent of E).

$$n_E(E) = g_{CB}(E) f(E) \quad , \quad g_{CB}(E) \propto \sqrt{E - E_c}$$

$$\Rightarrow g_{CB}(E) = A \sqrt{E - E_c}$$

↑
some constant

$$n_E(E) \approx A \sqrt{E - E_c} e^{-\frac{E - E_F}{k_B T}}$$

$$= A \sqrt{E - E_c} e^{-\frac{E}{k_B T}} e^{\frac{E_F}{k_B T}}$$

$$= A \sqrt{k_B T} \sqrt{\frac{E - E_c}{k_B T}} e^{-\frac{E}{k_B T}} e^{\frac{E_F}{k_B T}} e^{-\frac{E_c}{k_B T}} e^{\frac{E_c}{k_B T}}$$

$$= A \sqrt{k_B T} e^{-\frac{E_c - E_F}{k_B T}} \sqrt{\frac{E - E_c}{k_B T}} e^{-\frac{E - E_c}{k_B T}}$$

$$\Rightarrow n_x(x) = C\sqrt{x} e^{-x}$$

where $C = A\sqrt{k_B T} e^{-\frac{E_c - E_F}{k_B T}}$ and

$$x = \frac{E - E_c}{k_B T} \quad \square$$

b) Setting arbitrarily $C=1$, see the attached figure.

Maximum of the curve:

$$\frac{dn_x(x)}{dx} = \frac{1}{2\sqrt{x}} e^{-x} - \sqrt{x} e^{-x} = 0$$

$$e^{-x} \left(\frac{1}{2\sqrt{x}} - \sqrt{x} \right) = 0 \quad ||: e^{-x} \neq 0$$

$$\frac{1}{2\sqrt{x}} - \sqrt{x} = 0$$

$$\frac{1}{2\sqrt{x}} = \sqrt{x} \quad || \cdot \sqrt{x}$$

$$\underline{x = \frac{1}{2}}$$

$$n_{\max} = n_x\left(\frac{1}{2}\right) = \sqrt{\frac{1}{2}} e^{-\frac{1}{2}} \approx \underline{0,4289}$$

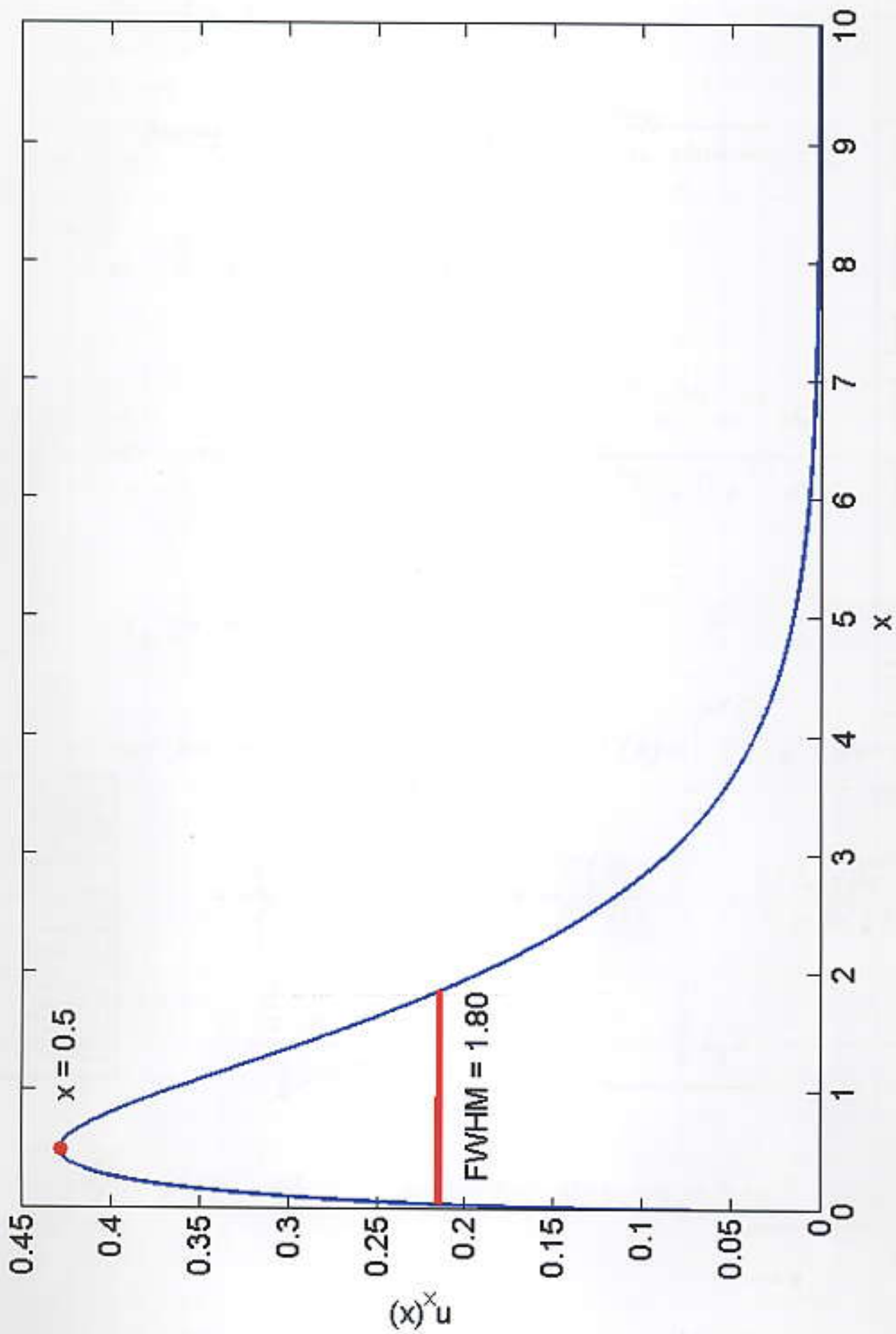
Full width at half maximum (FWHM) of the curve:

$$\sqrt{x} e^{-x} = \frac{n_{\max}}{2}$$

Numerically: $x_1 \approx 0,051$ and $x_2 \approx 1,846$

$$\Rightarrow \text{FWHM} = x_2 - x_1 \approx 1,795 \approx \underline{1,8}$$

$$x = \frac{E - E_c}{k_B T} \Rightarrow \text{energy from } E_c \text{ in units of } k_B T.$$



c) Let's show that the average electron energy in the CB is $\frac{3}{2} k_B T$.

$$\text{Average electron energy} = \frac{\text{Total energy}}{\text{Number of electrons}}$$

$$\Rightarrow \hat{E} = \frac{\int_{E_c}^{E_c+x} E n_E(E) dE}{\int_{E_c}^{E_c+x} n_E(E) dE}$$

$$\Rightarrow \hat{x} = \frac{\int_0^x x \sqrt{x} e^{-x} dx}{\int_0^x \sqrt{x} e^{-x} dx} = \frac{\int_0^x x^{3/2} e^{-x} dx}{\int_0^x x^{1/2} e^{-x} dx}$$

$$x \rightarrow \infty$$

$$\Rightarrow \hat{x} = \frac{\int_0^{\infty} x^{3/2} e^{-x} dx}{\int_0^{\infty} x^{1/2} e^{-x} dx}$$

$$\Gamma(z) = \int_0^{\infty} t^{z-1} e^{-t} dt$$

$$= \frac{\int_0^{\infty} x^{5/2-1} e^{-x} dx}{\int_0^{\infty} x^{3/2-1} e^{-x} dx} = \frac{\Gamma(5/2)}{\Gamma(3/2)}$$

$$\Gamma(1/2) = \sqrt{\pi} \quad \Gamma(z+1) = z \Gamma(z)$$

$$= \frac{\frac{3}{2} \cdot \frac{1}{2} \Gamma(1/2)}{\frac{1}{2} \Gamma(1/2)} = \frac{3}{2} \Rightarrow \hat{E} = \frac{3}{2} k_B T$$

$$\begin{aligned} \Gamma\left(\frac{5}{2}\right) &= \Gamma\left(\frac{3}{2}+1\right) = \frac{3}{2} \Gamma\left(\frac{3}{2}\right) \\ &= \frac{3}{2} \Gamma\left(\frac{1}{2}+1\right) = \frac{3}{2} \cdot \frac{1}{2} \Gamma\left(\frac{1}{2}\right) \end{aligned}$$

X/B! Average wavelength is calculated in a similar way from spectrum

d) $x = \frac{1}{2} \Rightarrow \frac{1}{2} = \frac{E - E_c}{k_B T} \Rightarrow E = E_c + \frac{1}{2} k_B T$
 $\underbrace{\hspace{10em}}_{E_{max}}$

Extrinsic n-Si

- 2) Si crystal has been doped n-type with $1 \cdot 10^{17} \text{ cm}^{-3}$ phosphorus (P) donors. The electron drift mobility μ_e depends on the total concentration of ionized dopants N_{dopant} .

Conductivity of the crystal:

Dopant concentration $N_d = 10^{17} \text{ cm}^{-3}$:

Drift mobility in Si for electrons $\mu_e = 730 \text{ cm}^2 \text{ V}^{-1} \text{ s}^{-1}$ and for holes $\mu_h = 328 \text{ cm}^2 \text{ V}^{-1} \text{ s}^{-1}$.

n-type conductivity:

Assuming $n = N_d$ and $p = \frac{n_i^2}{N_d}$ the conductivity becomes

$$\sigma = e N_d \mu_e + e \frac{n_i^2}{N_d} \mu_h \approx e N_d \mu_e$$

$$e = 1.602 \cdot 10^{-19} \text{ C}$$

$$n_i = 1.45 \cdot 10^{10} \text{ cm}^{-3}$$

$$\Rightarrow \sigma \approx 11.67 \text{ C cm}^{-1} \text{ V}^{-1} \text{ s}^{-1} = \underline{\underline{11.67 \frac{1}{\Omega \text{ cm}}}}$$

Fermi level with respect to the intrinsic crystal:

Intrinsic Fermi level:

$$E_{Fi} = E_v + \frac{1}{2} E_g - \frac{1}{2} k_B T \ln \frac{N_c}{N_v}$$

For notice, $E_{Fn} - E_{Fi}$ not calculated using this

$$n_i = N_c e^{-\frac{E_c - E_{Fi}}{k_B T}} \quad \text{e concentration in CB for intrinsic Si}$$

$$N_d = n = N_c e^{-\frac{E_c - E_{Fn}}{k_B T}} \quad \text{e concentration in CB for doped Si}$$

$$\Rightarrow \frac{N_d}{n_i} = \frac{N_c e^{-\frac{E_c - E_{Fn}}{k_B T}}}{N_c e^{-\frac{E_c - E_{Fi}}{k_B T}}} = e^{\frac{-E_c + E_{Fn} + E_c - E_{Fi}}{k_B T}} = e^{\frac{E_{Fn} - E_{Fi}}{k_B T}}$$

$$\Rightarrow \frac{E_{Fn} - E_{Fi}}{k_B T} = \ln \frac{N_d}{n_i} \Rightarrow E_{Fn} - E_{Fi} = k_B T \ln \frac{N_d}{n_i} \approx \underline{\underline{0.4 \text{ eV}}}$$

at room temp.

above intrinsic Fermi level

3) Compensation doping in n-type Si

N-type Si sample has been doped with 10^{16} phosphorus (P) atoms per cm^3

a) Electron and hole concentration: n-type

$$\text{Electron conc. : } n \approx N_d = \underline{\underline{10^{16} \text{ cm}^{-3}}}$$

$$\text{Hole conc. : } p = \frac{n_i^2}{N_d} = \frac{(1,45 \cdot 10^{10} \text{ cm}^{-3})^2}{10^{16} \text{ cm}^{-3}} \approx \underline{\underline{21000 \text{ cm}^{-3}}}$$

b) Room temperature conductivity of the sample:

$$\sigma = e N_d \mu_e + e \frac{n_i^2}{N_d} \mu_h$$

$$e = 1,602 \cdot 10^{-19} \text{ C}, \mu_e = 1200 \text{ cm}^2 \text{ V}^{-1} \text{ s}^{-1}, \mu_h = 444 \text{ cm}^2 \text{ V}^{-1} \text{ s}^{-1}$$

$$\Rightarrow \sigma = 1,92 \Omega^{-1} \text{ cm}^{-1}$$

c) Fermi level with respect to E_{Fi} :

$$n_i = N_c e^{-\frac{E_c - E_{Fi}}{k_B T}}$$

$$N_d = n = N_c e^{-\frac{E_c - E_{Fn}}{k_B T}}$$

$$\frac{N_d}{n_i} = \frac{N_c e^{-\frac{E_c - E_{Fn}}{k_B T}}}{N_c e^{-\frac{E_c - E_{Fi}}{k_B T}}} = e^{\frac{E_{Fn} - E_{Fi}}{k_B T}}$$

$$\Rightarrow E_{Fn} - E_{Fi} = k_B T \ln \frac{N_d}{n_i} \approx \overset{\text{at room temp.}}{\underline{\underline{0,35 \text{ eV}}}} \text{ above intrinsic Fermi level}$$

d) The crystal is now doped with 10^{17} boron acceptors per cm^3 . Electron and hole concentrations:

$$N_a = 10^{17} \text{ cm}^{-3} > N_d$$

$$\Rightarrow \text{Hole conc.: } p = N_a - N_d = \underline{9 \cdot 10^{16} \text{ cm}^{-3}}$$

$$\text{Electron conc.: } n = \frac{n_i^2}{p} \approx \underline{2740 \text{ cm}^{-3}}$$

e) Fermi level with respect to E_{Fi} :

Intrinsic:

Doped:

$$p_i = n_i \approx N_v e^{-\frac{E_{Fi} - E_v}{k_B T}} \quad , \quad p = N_a = N_v e^{-\frac{E_{Fp} - E_v}{k_B T}}$$

$$\Rightarrow \frac{p}{n_i} = \frac{N_v e^{-\frac{E_{Fp} - E_v}{k_B T}}}{N_v e^{-\frac{E_{Fi} - E_v}{k_B T}}} = e^{-\frac{E_{Fp} - E_{Fi}}{k_B T}}$$

$$\Rightarrow E_{Fp} - E_{Fi} = -k_B T \ln \frac{p}{n_i} \approx \underset{\substack{\text{at room} \\ \text{temp.}}}{-0.81 \text{ eV}} \quad \begin{array}{l} \text{below intrinsic} \\ \text{Fermi level} \end{array}$$

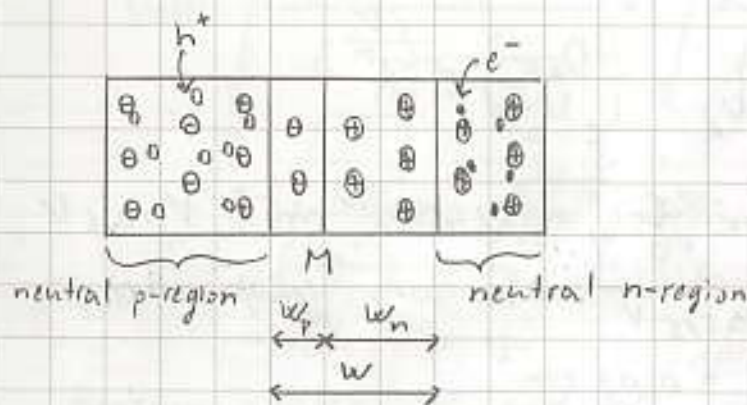
4) Si pn junction

Let's consider a long pn junction diode with an acceptor doping, N_a , of 10^{15} cm^{-3} on the p-side and donor concentration, N_d , on the n-side. The diode is forward biased with a voltage of 0.6V across it. The diode cross-sectional area is 1 mm^2 . The minority carrier recombination time, τ , depends on the dopant concentration, $N_{\text{dopant}} (\text{cm}^{-3})$, through the following approximate empirical relation $\tau \approx (5 \cdot 10^{-7}) / (1 + 2 \cdot 10^{13} N_{\text{dopant}})$ in seconds.

a) Let's suppose that $N_d = 10^{15} \text{ cm}^{-3}$. The depletion layer extends essentially into the n-side and we need to consider minority carrier recombination time, τ_n , in this region.

Let's calculate the diffusion and recombination contributions to the total diode current with $N_a = 10^{18} \text{ cm}^{-3}$.

Since $N_a \gg N_d$ this diode is p-n diode and $w_n > w_p$



Hole lifetime τ_h in the n-side: $N_{\text{dopant}} = 10^{15}$

$$\tau_h = \frac{5 \cdot 10^{-7}}{1 + 2 \cdot 10^{-17} \cdot 10^{15}} \text{ s} \approx 490,2 \text{ ns}$$

Electron lifetime τ_e in the p-side: $N_{\text{dopant}} = 10^{18}$

$$\tau_e = \frac{5 \cdot 10^{-7}}{1 + 2 \cdot 10^{-17} \cdot 10^{18}} \text{ s} \approx 23,8 \text{ ns}$$

Diffusion component of the diode current:

Diffusion coefficients of holes and electrons:

$$\mu_h = 478 \text{ cm}^2 \text{ V}^{-1} \text{ s}^{-1}, \mu_e = 280 \text{ cm}^2 \text{ V}^{-1} \text{ s}^{-1}$$

$$D_h = \frac{k_B T}{e} \mu_h \approx 12,356 \text{ cm}^2 \text{ s}^{-1}$$

$$D_e = \frac{k_B T}{e} \mu_e \approx 7,238 \text{ cm}^2 \text{ s}^{-1}$$

Diffusion lengths:

$$\text{into n-side: } L_h = \sqrt{D_h \tau_h} \approx 2,46 \cdot 10^{-3} \text{ cm} = 24,6 \text{ } \mu\text{m}$$

$$\text{into p-side: } L_e = \sqrt{D_e \tau_e} \approx 0,41 \cdot 10^{-3} \text{ cm} = 4,1 \text{ } \mu\text{m}$$

Thus diffusion component of the diode current is:

$$I_0 = A J_0 = A e n_i^2 \left(\frac{D_n}{L_n N_d} + \frac{D_p}{L_p N_a} \right) \underbrace{\left(e^{\frac{eV}{k_B T}} - 1 \right)}_{\approx e^{\frac{eV}{k_B T}}}$$

$$\approx A e n_i^2 \left(\frac{D_n}{L_n N_d} + \frac{D_p}{L_p N_a} \right) e^{\frac{eV}{k_B T}}$$

$$e = 1.602 \cdot 10^{-19} \text{ C}, \quad n_i = 1.45 \cdot 10^{10} \text{ cm}^{-3}, \quad V = 0.6 \text{ V},$$

$$\frac{e}{k_B T} = \frac{1}{0.02585 \text{ V}} \quad \text{at room temperature,}$$

$$A = 1 \text{ mm}^2 = 0.01 \text{ cm}^2$$

$$\Rightarrow I \approx 0.02042 \text{ A} = \underline{20.42 \text{ mA}}$$

$$\text{Unit: } \frac{\text{cm}^2 \cdot (\text{cm}^{-3})^2 \cdot \frac{\text{cm}^2}{\text{cm} \cdot \text{cm}^{-3}}}{\text{cm} \cdot \text{cm}^{-3}}$$

$$= \text{cm}^2 \cdot \text{cm}^{-6} \cdot \frac{\text{cm}^2}{\text{cm}^{-2} \text{ cm}}$$

$$= \frac{\text{cm}^2 \cdot \text{cm}^{-6} \cdot \text{cm}^4}{\text{cm}} = \frac{\text{C}}{\text{s}} = \text{A}$$

Recombination component of the diode current:

Built in potential:

$$V_0 = \frac{k_B T}{e} \ln \left(\frac{N_a N_d}{n_i^2} \right) \overset{\text{at room temp.}}{\approx} 0.75 \text{ V}$$

Knowing built in potential we are able to calculate the total width of the depletion region.

$$V_0 - V = \frac{e N_a N_d W_0^2}{2 \epsilon (N_a + N_d)} \Rightarrow W_0 = \sqrt{\frac{2 \epsilon (N_a + N_d) (V_0 - V)}{e N_a N_d}}$$

where $\epsilon = \epsilon_0 \epsilon_r = 1.034 \cdot 10^{-10} \text{ Fm}^{-1} = 1.034 \cdot 10^{-12} \text{ Fcm}^{-1}$
for silicon Si.

$$\Rightarrow w_0 \approx 4,47 \cdot 10^{-5} \text{ cm} = 0,447 \mu\text{m}$$

$$\begin{aligned} \text{Unit: } & \sqrt{\frac{\frac{\text{F}}{\text{cm}} \cdot \frac{1}{\text{cm}^3} \cdot \text{V}}{\text{C} \cdot \frac{1}{\text{cm}^3} \cdot \frac{1}{\text{cm}^3}}} = \sqrt{\frac{\frac{\text{C}}{\text{cm}} \cdot \frac{1}{\text{cm}^3} \cdot \cancel{\text{V}}}{\text{C} \cdot \frac{1}{\text{cm}^6}}} \\ & = \sqrt{\frac{\frac{1}{\text{cm}^4}}{\frac{1}{\text{cm}^6}}} = \sqrt{\text{cm}^2} = \text{cm} \end{aligned}$$

Depletion widths, w_p and w_n , can be calculated with N_a and N_d .

$$N_a w_p = N_d w_n, \quad w_p + w_n = w_0 \Rightarrow w_n = w_0 - w_p$$

$$N_a w_p = N_d (w_0 - w_p)$$

$$w_p (N_a + N_d) = N_d w_0$$

$$w_p = \frac{N_d w_0}{N_a + N_d} \approx 4,46 \cdot 10^{-5} \text{ cm} = 0,446 \mu\text{m}$$

$$w_n = w_0 - w_p \approx 4,47 \cdot 10^{-5} \text{ cm} = 0,447 \mu\text{m}$$

Thus the recombination component of the diode current is:

$$\begin{aligned} I_R = A J_R &= A \frac{en_i}{2} \left[\frac{w_p}{\tau_e} + \frac{w_n}{\tau_h} \right] \left[e^{\frac{eV}{2k_B T}} - 1 \right] \\ &\approx \frac{A en_i}{2} \left[\frac{w_p}{\tau_e} + \frac{w_n}{\tau_h} \right] e^{\frac{eV}{2k_B T}} \approx 0,00012 \text{ A} = \underline{\underline{0,12 \text{ mA}}} \end{aligned}$$

$$\text{Unit: } \text{cm}^2 \cdot \text{C} \cdot \frac{1}{\text{cm}^3} \cdot \frac{\cancel{\text{cm}}}{\text{s}} = \frac{\text{C}}{\text{s}} = \text{A}$$

Conclusion: Diode current is dominated by the diffusion component, $I_0 \gg I_R$.

b) Let's suppose that $N_d = N_a = 10^{18} \text{ cm}^{-3}$. Then w extends equally to both sides and further, $\tau_e = \tau_h$.

Let's calculate the diffusion and recombination contributions for this diode with the same procedure that in the previous case.

Hole and electron life time: $\tau_e = \tau_h \approx 23,8 \text{ ns}$

Hole and electron diffusion coefficients: $D_h \approx 12,756 \text{ cm}^2 \text{ s}^{-1}$
 $D_e \approx 7,238 \text{ cm}^2 \text{ s}^{-1}$

Diffusion lengths: $L_h = \sqrt{D_h \tau_h} \approx 0,54 \cdot 10^{-3} \text{ cm} = 5,4 \text{ } \mu\text{m}$

$L_e = \sqrt{D_e \tau_e} \approx 0,41 \cdot 10^{-3} \text{ cm} = 4,1 \text{ } \mu\text{m}$

Diffusion component of the diode current:

$$I_D = A J_D = A e n_i^2 \left(\frac{D_h}{L_h N_d} + \frac{D_e}{L_e N_a} \right) e^{\frac{eV}{k_B T}} \approx 1,62 \cdot 10^{-4} \text{ A} = \underline{\underline{0,162 \text{ mA}}}$$

Built in potential: $V_0 = \frac{k_B T}{e} \ln \left(\frac{N_a N_d}{n_i^2} \right) \approx 0,93 \text{ V}$ at room temp.

Total width of the depletion region:

$$w_0 = \sqrt{\frac{2 \epsilon (N_a + N_d) (V_0 - V)}{e N_a N_d}} \approx 2,93 \cdot 10^{-6} \text{ cm} = 29,3 \text{ nm}$$

Depletion widths, w_p and w_n : $w_p = w_n \approx 1,47 \cdot 10^{-6} \text{ cm} = 14,7 \text{ nm}$

Recombination component of the diode current:

$$I_R = A J_R = \frac{A e n_i}{2} \left[\frac{w_p}{\tau_e} + \frac{w_n}{\tau_h} \right] e^{\frac{eV}{2k_B T}} \approx 1,57 \cdot 10^{-4} \text{ A} = \underline{\underline{0,157 \text{ mA}}}$$

Conclusion: Diffusion and recombination components are nearly equal.